

A STUDY ON NANOFUID FLOW DUE TO A ROTATING DISK WITH SLIP CONDITION

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ABSTRACT

In this article, the effect of MHD viscous flow due to rotating disk is studied in the presence of heat generation/absorption and chemical reaction with velocity slip condition. The nonlinear ordinary differential equations are obtained and solved numerically using Runge-kutta fifth order algorithm . The effects of thermophysical parameters such as Prandtl number, Brownian motion, thermophoresis parameter on velocity, temperature and concentration are discussed in the help of graphs through MATLAB. Further the numerical values of HTR and MTR are offered with the aid of table. In addition the present outcomes obtained is compared with existing published work.

KEYWORDS: MHD, Heat generation/absorption, chemical reaction, partial slip condition.

INTRODUCTION

Magnetohydrodynamics denotes the study of the motion of highly electrically conducting fluids. The role of MHD is important in many fields like nuclear reactors, solar physics etc. Nanofluid is a new kind of heat transfer medium, containing Nanoparticles. It plays an important role in the enhancement of thermal conductivity and can be used for welding equipments, crystal silicon mirror cooling. Fluid flow evoked by a rotating disk has been a compelling analysis topic since it is relevant during a range of technical applications involving chemical science systems, deposition of coatings on surfaces, rotor-stator system, atmospheric and oceanic circulations, viscometer and varied others. Kumaran(2010) studied the transition of MHD boundary layer flow past a stretching sheet. He investigated the transition effects because of applied magnetic field were analysed numerically on Newtonian fluid flow. Whereas Noor N G(2010) carried out the simple non-perturbative solution of MHD Newtonian fluid flow over a shrinking surface. Turkyilmazoglu M(2012) studied the exact analytical solutions for heat and mass transfer of MHD slip flow in nanofluids. It investigates the influence of slip on the behavior of fluid flow and thermal transport of some electrically conducting nanofluid over a permeable stretching/shrinking sheet. In 2013 Turkyilmazoglu M he studied the heat and mass transfer flow due to rotating rough and porous disk the rotating disk surface is considered with partial slip in the presence of a uniform suction or injection. He analysed the Effects of wall roughness and temperature jump on the heat and mass transfer. In 2013 Dandapat and Singh has reviewed the unsteadiness effect in two layer film flow on rotating disc. Hayat, T (2015) have studied the magnetohydrodynamic (MHD) flow of viscous nanofluid saturating porous medium. He investigates that higher nanoparticle volume fraction decreases the velocity field also an analysis has been carried out for Magnetohydrodynamic three-dimensional flow of viscoelastic nanofluid in the presence of nonlinear. He analysed the MHD 3D flow of nanofluid in the presence of thermophoresis and Brownian motion effects. The experimental results shows that the thermal boundary layer thickness is an increasing function of radiative effect.

Therefore the objective of the present work is to analyse the effects of MHD viscous flow due to rotating disk in the presence of heat generation and chemical reaction. The present work is likely to have bearing on the problem of heat transfer which can be useful in industries for designing rotating machineries and lubrications etc. The objective of this article is to explore the study of viscous nanofluid flow due to rotating rigid disk with slip condition for MHD and hydrodynamic cases. The mutual interaction of thermophoretic and Brownian motion phenomena is acknowledged by considering nano particles. The effects of various thermophysical

parameters like Magnetic parameter B , Velocity slip parameter λ , Lewis number Le , Prandtl number Pr , Brownian motion parameter N_B and thermophoresis parameter N_T on the velocity, temperature and concentration profiles are examined graphically through MATLAB.

MATHEMATICAL FORMULATION

The viscous Nanofluid flow brought by a rotating disk subject to velocity slip condition is considered. In the direction of \bar{z} $-axis$ the uniform magnetic field with strength B_0 is applied. At $\bar{z} = 0$, the disk rotates with a constant angular velocity $\bar{\Omega}$. The problem is take into account with low Reynolds number so that the induced magnetic fields is neglected. The effects of heat generation/absorption and chemical reaction are examined. The phenomena of mutual interaction of thermophoresis and brownian motion in heat and mass transfer aspects is considered. $(\bar{u}, \bar{v}, \bar{w})$ are the velocity components in the direction of $(\bar{r}, \bar{\phi}, \bar{z})$. The boundary layer approximation which reduces the governing equation to present flow situation as follows:

Continuity Equation

$$\frac{\partial \bar{w}}{\partial \bar{z}} + \frac{\partial \bar{u}}{\partial \bar{r}} + \frac{\bar{u}}{\bar{r}} = 0 \tag{1}$$

Momentum Equation

$$\bar{w} \frac{\partial \bar{u}}{\partial \bar{z}} + \bar{u} \frac{\partial \bar{u}}{\partial \bar{r}} - \frac{\bar{v}^2}{\bar{r}} = \nu \left(\frac{\partial^2 \bar{u}}{\partial \bar{z}^2} + \frac{1}{\bar{r}} \frac{\partial \bar{u}}{\partial \bar{r}} + \frac{\partial^2 \bar{u}}{\partial \bar{r}^2} - \frac{\bar{u}}{\bar{r}^2} \right) - \frac{\sigma B_0^2}{\rho_f} \bar{u} \tag{2}$$

$$\bar{w} \frac{\partial \bar{v}}{\partial \bar{z}} + \bar{u} \frac{\partial \bar{v}}{\partial \bar{r}} + \frac{\bar{w}\bar{v}}{\bar{r}} = \nu \left(\frac{\partial^2 \bar{v}}{\partial \bar{z}^2} + \frac{\partial^2 \bar{v}}{\partial \bar{r}^2} - \frac{\bar{v}}{\bar{r}^2} + \frac{1}{\bar{r}} \frac{\partial \bar{v}}{\partial \bar{r}} \right) - \frac{\sigma B_0^2}{\rho_f} \bar{v} \tag{3}$$

$$\bar{w} \frac{\partial \bar{w}}{\partial \bar{z}} + \bar{u} \frac{\partial \bar{w}}{\partial \bar{r}} = \nu \left(\frac{\partial^2 \bar{w}}{\partial \bar{z}^2} + \frac{1}{\bar{r}} \frac{\partial \bar{w}}{\partial \bar{r}} + \frac{\partial^2 \bar{w}}{\partial \bar{r}^2} \right) \tag{4}$$

Heat Equation

$$\bar{w} \frac{\partial \bar{T}}{\partial \bar{z}} + \bar{u} \frac{\partial \bar{T}}{\partial \bar{r}} = \alpha \left(\frac{\partial^2 \bar{T}}{\partial \bar{z}^2} + \frac{1}{\bar{r}} \frac{\partial \bar{T}}{\partial \bar{r}} + \frac{\partial^2 \bar{T}}{\partial \bar{r}^2} \right) + \frac{(\rho c)_p}{(\rho c)_f} \left[D_B \left(\frac{\partial \bar{T}}{\partial \bar{z}} \frac{\partial \bar{C}}{\partial \bar{z}} + \frac{\partial \bar{T}}{\partial \bar{r}} \frac{\partial \bar{C}}{\partial \bar{r}} \right) \right] + \frac{(\rho c)_p}{(\rho c)_f} \left[\frac{D_T}{\bar{T}_\infty} \left(\left(\frac{\partial \bar{T}}{\partial \bar{z}} \right)^2 + \left(\frac{\partial \bar{T}}{\partial \bar{r}} \right)^2 \right) \right] + \frac{Q_0}{c_{p\rho}} (\bar{T} - \bar{T}_\infty) \tag{5}$$

Mass Equation

$$\bar{w} \frac{\partial \bar{C}}{\partial \bar{z}} + \bar{u} \frac{\partial \bar{C}}{\partial \bar{r}} = D_B \left(\frac{\partial^2 \bar{C}}{\partial \bar{z}^2} + \frac{1}{\bar{r}} \frac{\partial \bar{C}}{\partial \bar{r}} + \frac{\partial^2 \bar{C}}{\partial \bar{r}^2} \right) + \frac{D_T}{\bar{T}_\infty} \left[\frac{\partial^2 \bar{T}}{\partial \bar{z}^2} + \frac{1}{\bar{r}} \frac{\partial \bar{T}}{\partial \bar{r}} + \frac{\partial^2 \bar{T}}{\partial \bar{r}^2} \right] - R_0 (\bar{C} - \bar{C}_\infty) \tag{6}$$

The corresponding boundary conditions are:

$$\bar{u} = L \frac{\partial \bar{u}}{\partial \bar{z}}, \bar{v} = \bar{r}\bar{\Omega} + L \frac{\partial \bar{v}}{\partial \bar{z}}, \bar{w} = 0, \bar{T} = \bar{T}_W, \bar{C} = \bar{C}_W \text{ as } \bar{z} = 0 \tag{7}$$

$$\bar{u} \rightarrow 0, \bar{v} \rightarrow 0, \bar{T} \rightarrow \bar{T}_\infty, \bar{C} \rightarrow \bar{C}_\infty \text{ as } \bar{z} \rightarrow \infty. \tag{8}$$

where, $\nu = \frac{\mu}{\rho_f}$ is the kinematic viscosity, μ is the dynamic viscosity, σ is the electrical conductivity of the fluid, ρ is the density, ρ_f is the density of the base fluid, $(\rho c)_p$ is the effective heat capacity of nanoparticles, $\alpha = \frac{k}{(\rho c)_f}$ is the thermal diffusivity, $(\rho c)_f$ is the heat capacity, \bar{T} is temperature, \bar{C} is the concentration, D_T is the thermophoretic diffusion coefficient, D_B is the Brownian diffusion coefficient, L is the velocity slip constant, \bar{T}_W is the surface temperature, \bar{T}_∞ is the ambient temperature, \bar{C}_W is the surface concentration, \bar{C}_∞ is the ambient concentration, Q_0 is the heat generation/absorption coefficient and R_0 is the rate of chemical reaction. The dimensionless variables that we introduced here is as follow

$$\bar{u} = \bar{r}\bar{\Omega} \frac{dF(\zeta)}{d\zeta}, \bar{v} = \bar{r}\bar{\Omega} G(\zeta), \bar{w} = -\sqrt{2\bar{\Omega}\nu} F(\zeta) \tag{9}$$

$$C(\zeta) = \frac{\bar{C} - \bar{C}_\infty}{\bar{C}_W - \bar{C}_\infty}, T(\zeta) = \frac{\bar{T} - \bar{T}_\infty}{\bar{T}_W - \bar{T}_\infty}, \zeta = \sqrt{\frac{2\bar{\Omega}}{\nu}} \bar{z}$$

By incorporating the equation (9), The equation (1) – (8) is transformed to the non linear ordinary differential equation follows: The equation (1) satisfied.

The Momentum equation (2) transforms to the ordinary differential equation as follows

$$2 \frac{d^3 F(\zeta)}{d\zeta^3} + 2F(\zeta) \frac{d^2 F(\zeta)}{d\zeta^2} - \left(\frac{dF(\zeta)}{d\zeta}\right)^2 + (G(\zeta))^2 - \beta \frac{dF(\zeta)}{d\zeta} = 0 \tag{10}$$

The Momentum equation (3) transforms to the ordinary differential equation as follows

$$2 \frac{d^2 G(\zeta)}{d\zeta^2} + 2F(\zeta) \frac{dG(\zeta)}{d\zeta} - 2G(\zeta) \frac{dF(\zeta)}{d\zeta} - \beta G(\zeta) = 0 \tag{11}$$

The Heat equation (4) transforms to the ordinary differential equation as follows

$$\frac{d^2 T(\zeta)}{d\zeta^2} + Pr \left(F(\zeta) \frac{dT(\zeta)}{d\zeta} + N_B \frac{dT(\zeta)}{d\zeta} \frac{dC(\zeta)}{d\zeta} + N_T \left(\frac{dT(\zeta)}{d\zeta}\right)^2 + HT(\zeta) \right) = 0 \tag{12}$$

The Mass equation (5) transforms to the ordinary differential equation as follows

$$\frac{d^2 C(\zeta)}{d\zeta^2} + LePrF(\zeta) \frac{dC(\zeta)}{d\zeta} + \frac{N_T}{N_B} \frac{d^2 T(\zeta)}{d\zeta^2} - R_p C(\zeta) = 0 \tag{13}$$

The boundary conditions (7) and (8) changes to

$$F(\zeta) = w_s, \frac{dF(\zeta)}{d\zeta} = \lambda \frac{d^2 F(\zeta)}{d\zeta^2}, G(\zeta) = 1 + \lambda \frac{dG(\zeta)}{d\zeta}, T(\zeta) = 1, C(\zeta) = 1 \text{ at } \zeta = 0 \tag{14}$$

$$\frac{dF(\zeta)}{d\zeta} \rightarrow 0, G(\zeta) \rightarrow 0, T(\zeta) \rightarrow 0, C(\zeta) \rightarrow 0 \text{ as } \zeta \rightarrow \infty$$

Here, β is the magnetic field parameter, λ is the velocity slip parameter, Pr is the Prandtl number, N_T is the thermophoresis parameter, N_B is the Brownian motion parameter, H is the heat generation/absorption parameter, R_p is the chemical reaction parameter, and Le is the Lewis number and these parameters are defined as follows:

$$\beta = \sqrt{\frac{\sigma \beta_0^2}{\rho_f \Omega}}, \lambda = L \sqrt{\frac{2\Omega}{\nu}}, N_B = \frac{(\rho C)_p}{(\rho C)_f} \frac{(\bar{T}_w - \bar{T}_\infty) D_T}{\bar{T}_\infty}, H = \frac{Q_0}{2\Omega \rho C_p}, Le = \frac{\alpha}{D_B}, Pr = \frac{\nu}{\alpha}, N_T = \frac{(\rho C)_p}{(\rho C)_f} \frac{(\bar{C}_w - \bar{C}_\infty) D_B}{\nu}, R_p = \frac{\nu}{2\Omega D_B} R_0 \tag{15}$$

the dimensionless forms of skin friction coefficient (SFC), local nusselt number (and local Sherwood number) are defined as

$$\sqrt{Re_{\bar{r}}} C_F = \frac{d^2 F(0)}{d\zeta^2}, \sqrt{Re_{\bar{r}}} C_G = \frac{dG(0)}{d\zeta}, \frac{Nu}{\sqrt{Re_{\bar{r}}}} = -\frac{dT(0)}{d\zeta}, \frac{Sh}{\sqrt{Re_{\bar{r}}}} = -\frac{dC(0)}{d\zeta} \tag{16} \text{ where } Re_{\bar{r}} = \frac{(\bar{\Omega} \bar{r}) \bar{r}}{2\nu} \text{ denoted the local rotational Reynolds number.}$$

NUMERICAL APPROACH

Mathematical equations are obtained in terms of higher order differential equation. Reduction of partial differential equation to non linear ordinary differential equation by using some suitable similarity transformation.

To solve the equations (10) - (13) numerically which is subjected to the boundary conditions (14), is executed through Runge-kutta fifth order algorithm. The system of equations is converted to initial value problem.

The substitutions used here are:

$$q_1 = F(\zeta), q_2 = F'(\zeta), q_3 = F''(\zeta), q_4 = G(\zeta), q_5 = G'(\zeta), q_6 = T(\zeta), q_7 = T'(\zeta), q_8 = C(\zeta), q_9 = C'(\zeta)$$

and obtained the following system of ordinary differential equations

$$\begin{aligned}
 q_1' &= q_2 \\
 q_2' &= q_3 \\
 q_3' &= \frac{(q_2)^2 - 2q_1q_3 - (q_4)^2 + \beta q_2}{2} \\
 q_4' &= q_5 \\
 q_5' &= q_2q_4 - q_1q_5 + \frac{1}{2}\beta^2q_4 \\
 q_6' &= q_7 \\
 q_7' &= -Pr[q_1q_7 + N_Bq_7q_9 + N_T(q_7)^2 + Hq_6] \\
 q_8' &= q_9 \\
 q_9' &= -LePrq_9 + \frac{N_T}{N_B}q_7' + R_Pq_8
 \end{aligned} \tag{17}$$

With conditions

$$q_1(0) = 0, q_2(0) = \lambda F''(0), q_3(0) = F''(0), q_4(0) = 1 + \lambda G'(0), q_5(0) = G'(0), q_6(0) = 1, q_8(0) = 1. \tag{18}$$

With additional conditions

$$q_2(\infty) = 0, q_4(\infty) = 0, q_6(\infty) = 0, q_8(\infty) = 0. \tag{19}$$

For solving the above system numerically, we implement the Runge-kutta fifth order algorithm. For the integration purpose of (17) as initial value problem we need 9 initial conditions but only 7 are given. Before starting solving the system of ode favourable guessed values for $q_7(0), q_9(0)$ are needed. But in addition we have the condition (19) so the upper limit of ζ is finite. Let 8 be the randomly chosen value for ζ_∞ and By considering suitable guesses for $q_7(0), q_9(0)$ as -1. The integration of the above system of first order non linear differential equations can carried out with the above assumptions.. The exact values of these initial guesses are computed through Newton Raphson method. In the course of computations, we have chosen h as 0.01 while residual of boundary conditions at infinity is 10^{-5} .

Tables 1 are constructed to offered numerical values of local nusselt number ($-T'(0)$) and local Sherwood number ($-C'(0)$) for different values of various parameters $\beta, \lambda, N_T, N_B, Le, Pr$.

Table 1. Numerical values of HTR ($-T'(0)$) and MTR ($-C'(0)$)

| λ | β | Le | Pr | N_B | N_T | $(-T'(0))$ | | $(-C'(0))$ | |
|-----------|---------|------|------|-------|-------|----------------|---------------|----------------|---------------|
| | | | | | | Existing value | Present value | Existing value | Present value |
| 0.2 | 0.3 | 0.8 | 1.0 | 0.3 | 0.2 | 0.32660 | 0.32640 | 0.27591 | 0.27564 |
| 0.5 | | | | | | 0.30360 | 0.30450 | 0.26942 | 0.26940 |
| 0.8 | | | | | | 0.28724 | 0.28781 | 0.26499 | 0.26309 |
| 0.7 | 0.0 | 0.8 | 1.0 | 0.3 | 0.2 | 0.30502 | 0.30464 | 0.27011 | 0.27039 |
| | 0.7 | | | | | 0.24434 | 0.24453 | 0.25395 | 0.25382 |
| | 1.4 | | | | | 0.17572 | 0.17545 | 0.23734 | 0.23480 |

| | | | | | | | | | |
|-----|-----|-----|-----|-----|-----|---------|---------|---------|---------|
| 0.7 | 0.3 | 0.5 | 1.0 | 0.3 | 0.2 | 0.29641 | 0.29663 | 0.21382 | 0.21360 |
| | | 1.0 | | | | 0.28961 | 0.28945 | 0.30142 | 0.30129 |
| | | 1.5 | | | | 0.28399 | 0.28315 | 0.38695 | 0.38890 |
| 0.7 | 0.3 | 0.8 | 0.5 | 0.3 | 0.2 | 0.24998 | 0.24985 | 0.22942 | 0.22939 |
| | | | 1.0 | | | 0.29233 | 0.29218 | 0.26632 | 0.26629 |
| | | | 1.5 | | | 0.32293 | 0.32285 | 0.31278 | 0.31258 |
| 0.7 | 0.3 | 0.8 | 1.0 | 0.5 | 0.2 | 0.26359 | 0.26359 | 0.30347 | 0.30339 |
| | | | | 0.7 | | 0.23688 | 0.23678 | 0.31882 | 0.31869 |
| | | | | 1.0 | | 0.20067 | 0.20051 | 0.32971 | 0.32963 |
| 0.7 | 0.3 | 0.8 | 1.0 | 0.3 | 0.5 | 0.25915 | 0.25905 | 0.22212 | 0.22202 |
| | | | | | 0.7 | 0.23878 | 0.23865 | 0.22561 | 0.22517 |
| | | | | | 1.0 | 0.21024 | 0.21020 | 0.22286 | 0.22272 |

FIGURES

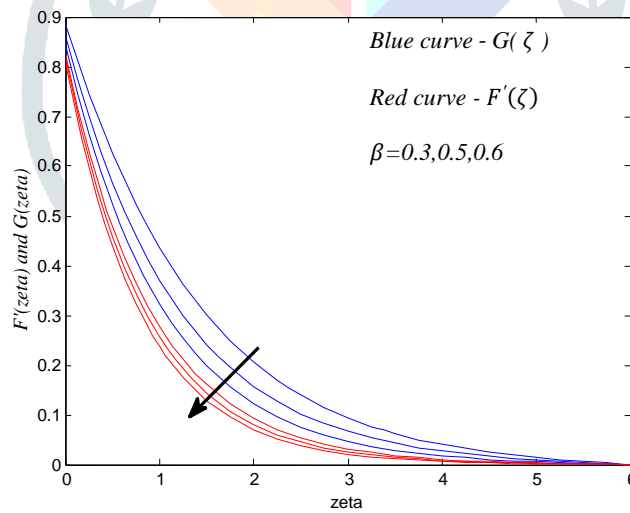


Figure 1 Profiles of $F'(\zeta)$ and $G(\zeta)$ for β

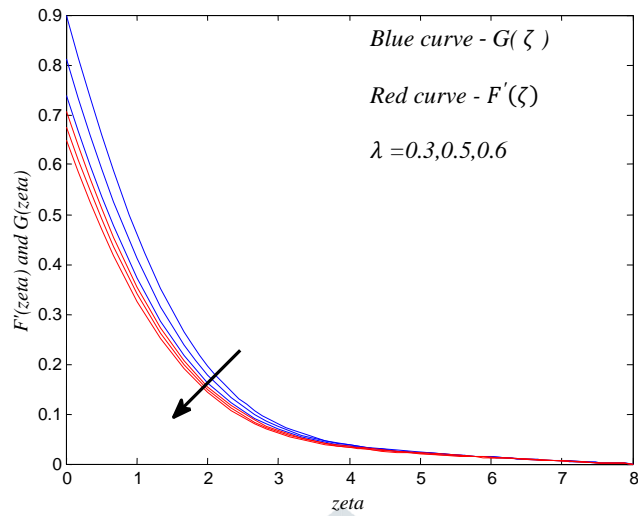


Figure 2 . Profiles of $F'(\zeta)$ and $G(\zeta)$ for λ

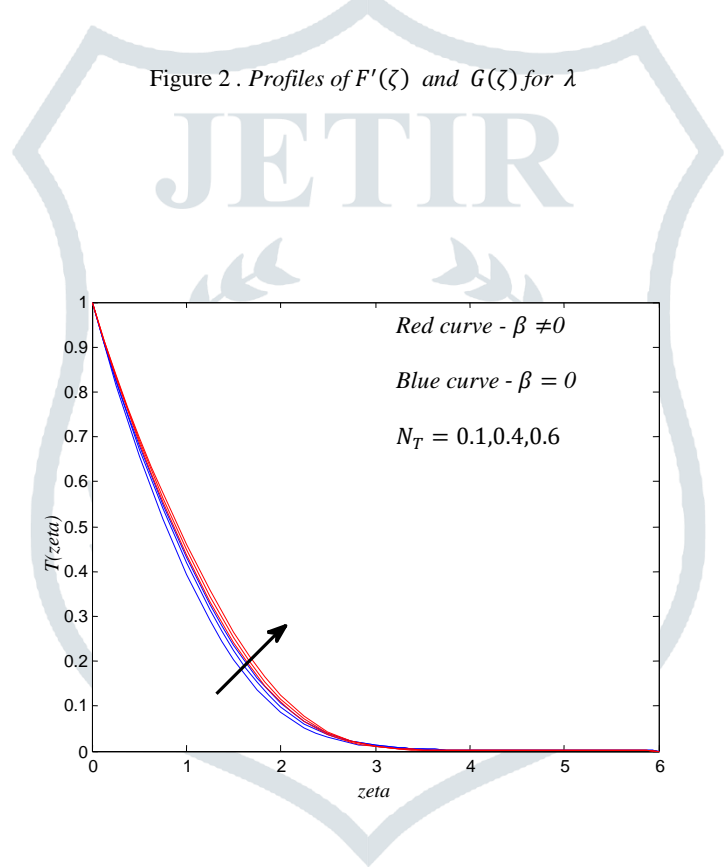


Figure 3. Profile of $T(\zeta)$ for N_T

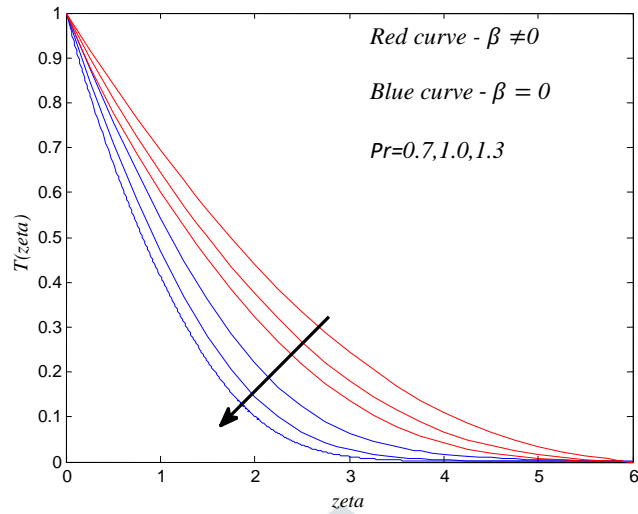


Figure 4. Profile of $T(\zeta)$ for Pr

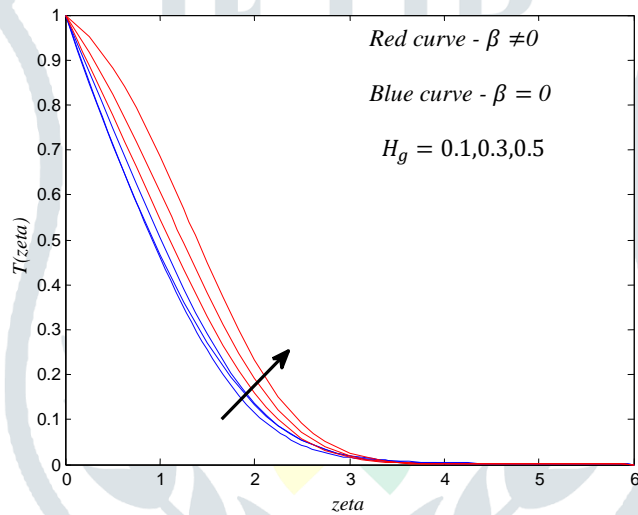


Figure 5. The Profile of $T(\zeta)$ for H_g

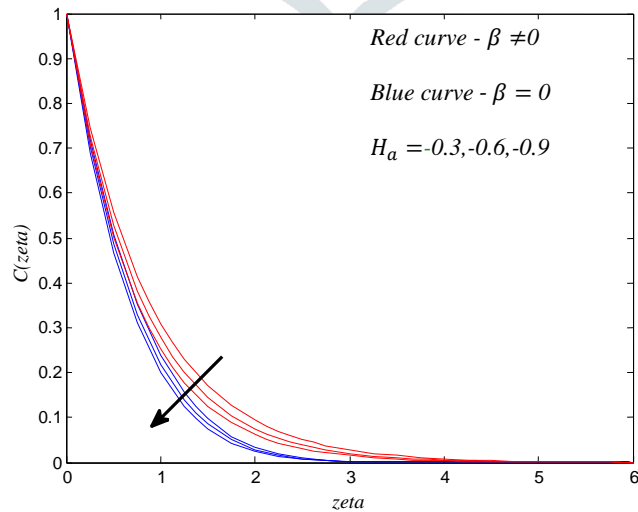


Figure 6. Profile of H_a on $T(\zeta)$

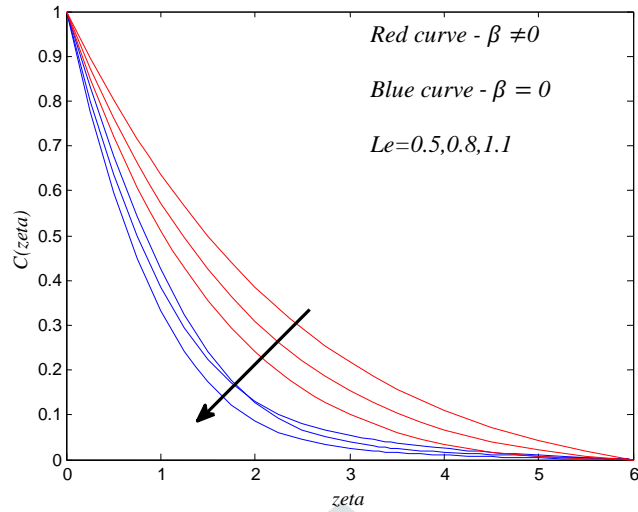


Figure 7. Profile of $C(\zeta)$ for Le

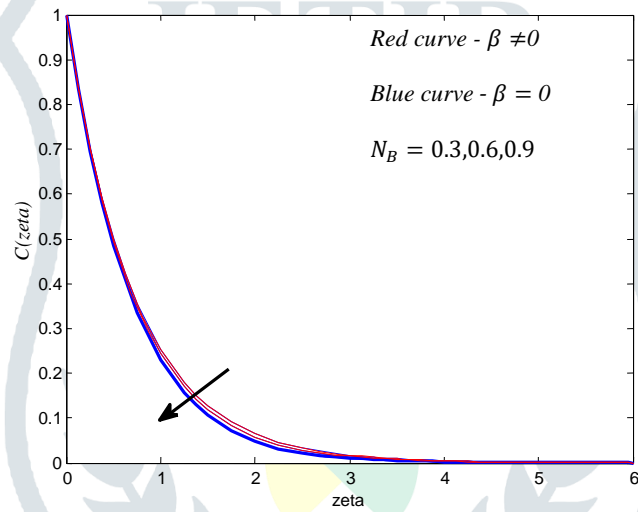


Figure 8. Profile of $C(\zeta)$ for N_B

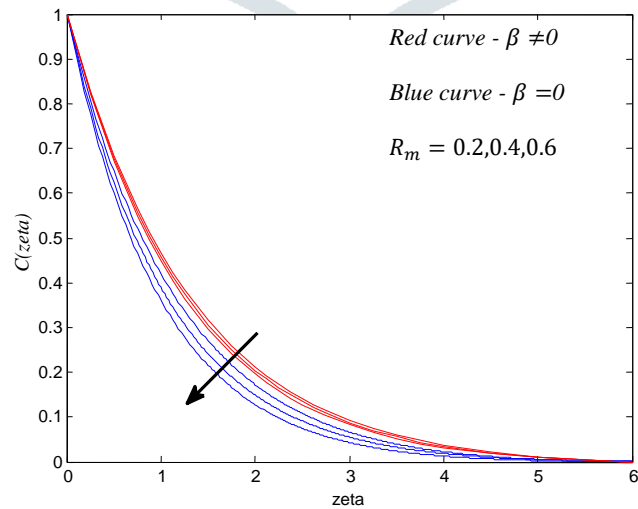


Figure 9. Profile of $C(\zeta)$ for R_m

RESULT AND DISCUSSION

A viscous nanofluid flow brought by a rotating disk subject to velocity slip condition is considered. Heat generation/absorption effect is taken into account by temperature equation while the chemical reaction effect is admitted by concentration equation. The properties of Thermophoresis and Brownian motion taking into account by Nanofluid model. Here in Table 1. We observed that Le and Pr on local nusselt number are quite opposite whereas gives the similar behavior for higher the values of Brownian motion and thermophoresis parameter. We also analysed that the local Sherwood number shows the increasing behavior for the larger values of Brownian motion and thermophoresis parameter.

Velocity Profiles

Figure 1. illustrates the viscous fluid velocities ($F'(\zeta)$ and $G(\zeta)$) for various values of magnetic field parameter. It can be seen that the increase in magnetic field parameter cause a decline effect on both $F'(\zeta)$ and $G(\zeta)$. Higher the value of β decreases the viscous fluid velocities. Figure 2. demonstrate the effect of λ on both $F'(\zeta)$ and $G(\zeta)$. It shows that the both velocities decreases by increasing the velocity slip parameter λ .

Temperature Profiles

In Figure 3. we examined the behavior of parameter N_T on temperature profile $T(\zeta)$ in the presence and absence of magnetic field parameter ($\beta \neq 0$ and $\beta = 0$). It shows that the increasing value N_T leads to higher temperature profile $T(\zeta)$. Figure 4. illustrates the effect of Prandtl number on temperature profile with absence/presence of β . It is noticed that fluid temperature is decreased via increasing value of N_T . Figure 5. shows the effects of heat generation on fluid temperature profile for both $\beta \neq 0$ and $\beta = 0$. It is seen that the $T(\zeta)$ is increasing function of H_g . Figure 6. illustrates the effects of heat absorption H_a of temperature profile on both zero and non zero values of magnetic field parameter β . It depicts that the larger value of H_a leads to a lower temperature profile.

Concentration Profiles

Figure 7-9. illustrates the influence of Le , N_B , R_m on nanoparticle concentration $C(\zeta)$. The impact of Le on concentration profile is examined in Fig.7. It is seen that for both zero and non zero values of β an increase in Le cause a decay in the concentration profile $C(\zeta)$. In Fig.8 the influence of Brownian motion parameter N_B on concentration profile $C(\zeta)$ is depicted. It is observed that the higher value of N_B brings decrease in fluid concentration for both $\beta \neq 0$ and $\beta = 0$. The remarkable changes in nanoparticle concentration is identified by the way of Fig.9 The positive value of chemical reaction parameter brings decline in fluid concentration $C(\zeta)$. More precisely for both $\beta \neq 0$ and $\beta = 0$ the concentration profiles shows an inciting nature for the higher values of R_m .

CONCLUSION

In this article both MHD and hydrodynamic frame of viscous nanofluid due to rotating disk with slip condition is examined. The fluid flow situation is narrated with partial slip condition, chemical reaction and heat sink/source. The numerical solution for solving governing non linear ordinary differential equation is executed Runge-kutta 5th order algorithm. The following conclusions are made from the present investigations.

It is observed that the fluid velocities $F'(\zeta)$ and $G(\zeta)$ are decreasing function of magnetic field parameter β and also it exhibits decline nature towards the higher value of slip parameter λ . The Fluid temperature $T(\zeta)$ is an increasing function for high values of thermophoresis parameter N_T whereas it is also observed that a decline nature towards Pr , H_g and H_a for both MHD and hydrodynamic case. The viscous Nanofluid concentration $C(\zeta)$ shows a decline attitude for Lewis number, Brownian motion and chemical reaction parameters in both the cases $\beta \neq 0$ and $\beta = 0$. It is observed that the fluid velocities $F'(\zeta)$ and $G(\zeta)$ are decreasing function of magnetic field parameter β and also it exhibits decline nature towards the higher value of slip parameter λ . The Fluid temperature $T(\zeta)$ is an increasing function for high values of thermophoresis parameter N_T . It is observed that a decline nature towards Pr , H_g and H_a for both MHD and hydrodynamic case. The viscous Nanofluid concentration $C(\zeta)$ shows a decline attitude for Lewis number, Brownian motion and chemical reaction parameters in both the cases $\beta \neq 0$ and $\beta = 0$.

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