

Estimation of effective Debye temperature of some liquids and their Isotopes by acoustic method

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Abstract: Acoustic methods have been applied to estimate the effective Debye temperature (θ_D) of pure liquids and their isotopes. The theoretical approach employed earlier has been completely modified and a new approach has been suggested. Five pure liquids and their isotopes have been undertaken for computing the values of effective Debye temperature using the present approach. The necessary experimental data needed have been taken from literature.

Key words: Debye temperature, ultrasonic velocity, temperature.

1. Introduction

The structure of a liquid can be denoted by the term quasi-crystalline in the sense that liquid displays a certain degree of local order of the same type as that of corresponding crystal. The study of dynamics of atomic motion of a liquid plays very important role in understanding the solid like behavior of liquids. So the name effective Debye temperature has been given. On the basis of quasi-crystalline structure and other experimental evidences, a number of workers¹⁻⁹ evaluated the effective Debye temperature of various liquids as a function of temperature and pressure. These workers utilized ultrasonic absorption, velocity, density and compressibility data of liquid for the estimation of effective Debye temperature. In the subsequent development⁸ the theoretical approach has been modified to enable the calculation of Debye temperature from sound velocity, density, compressibility and heat capacities ratio data. Qualitatively, the concept of effective Debye temperature has been introduced to few binary liquid mixtures in the year 1979¹⁰. Since then the study of effective Debye temperature has not been progressively done. In the present work, we are introducing a fresh approach to study the effective Debye temperature of pure liquids and their isotopes.

Previously, Debye temperature of liquids has been computed on the basis of approximate equivalence of β_T and β_a , just as in the case of solids. Here we have assumed $\beta_a \neq \beta_T$, as α has considerable values in liquids but not in solids. So $\gamma \neq 1$.

2. Theoretical

The expression for Debye temperature (θ_D) for a solid is given by

$$\theta_D = \frac{h}{k} \left[\left(\frac{9N}{4\pi V} \right) / \left(\frac{1}{C_l^3} + \frac{2}{C_t^3} \right) \right]^{1/3} \quad (1)$$

where h , k , N and V are Planck's constant, Boltzmann's constant, Avogadro number and molar volume respectively.

here

$$\left(\frac{1}{C_l^3} + \frac{2}{C_i^3}\right) = (\rho\beta_a)^{3/2} \left[\left\{ \frac{(1+\sigma)}{3(1-\sigma)} \right\}^{3/2} + 2 \left\{ \frac{2(1+\sigma)}{3(1-2\sigma)} \right\}^{3/2} \right] \quad (2)$$

In the above expression β_a can be written as follows

$$\beta_a = \beta_T - \frac{TV\alpha^2}{C_p} \quad (3)$$

where α , C_p , β_T are expansivity, the specific heat at constant pressure and the isothermal compressibility respectively. For the solid, α is very small, so the second term is neglected for solids, and,

$$\beta_a \approx \beta_T \quad (4)$$

Equation (4) has been used only for the solids, but recently some workers¹¹ used this assumption in calculating Debye temperature of liquids and liquid mixtures. In the case of liquids, α has considerable value so here we have modified above approach by assuming $\beta_a \neq \beta_T$ and β_a has been calculated from the relation

$$\beta_a = (\rho_{mix} u^2_{mix})^{-1} \quad (5)$$

Poisson ratio has been calculated as

$$\sigma = \frac{3A-2}{6A+2} \quad (6)$$

where

$$A = \left(\frac{K_{T,\infty}}{G_{T,\infty}} \right) = \left(\frac{K_{T,\infty}}{K_{T,r}} \right) \left(\frac{\eta_v}{\eta_s} \right) \left(\frac{\tau_s}{\tau_v} \right) \quad (7)$$

and using the relation given by Davis and Litovitz¹² we can write

$$\left(\frac{K_{T,\infty}}{K_{T,r}} \right) \left(\frac{\eta_v}{\eta_s} \right) \left(\frac{\tau_s}{\tau_v} \right) = \frac{4}{3\gamma}$$

$$A = \frac{4}{3\gamma} \quad (8)$$

where γ is defined as the ratio of isothermal compressibility to adiabatic compressibility given as

$$\gamma = \frac{\beta_T}{\beta_a} \quad (9)$$

where β_T is calculated from recently proposed relation¹³

$$\beta_T = \frac{1.71 \times 10^{-3}}{u^2 T^{4/9} \rho^{4/3}} \quad (10)$$

3. Results and Discussion

The values of Debye temperature of pure liquids viz. H₂O, C₆H₆, C₆H₁₂, CH₃OH, C₂H₅OH and their isotopes D₂O, C₆D₆, C₆D₁₂, CH₃OD, C₂H₅OD have been calculated at four different temperatures 293.15, 303.15, 313.15 and 323.15 K with the help of above procedures. The required data needed for the computations has been taken from literature¹⁴. Calculated values of Debye temperature of pure liquids are given in Table-1. Equations (1) to (10) have been used to obtain the values of Debye temperature of liquids and their isotopes under consideration.

A close perusal of table-1 shows that Debye temperature of pure liquids and their respective isotopes are decreasing with increase in temperature except H₂O and its isotope. This also indicates that Debye temperature decreases for heavy molecules. This is due to increase in density and decrease in ultrasonic velocity of heavy molecules.

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Table1: Debye Temperature of pure liquids and its isotopes at different temperatures

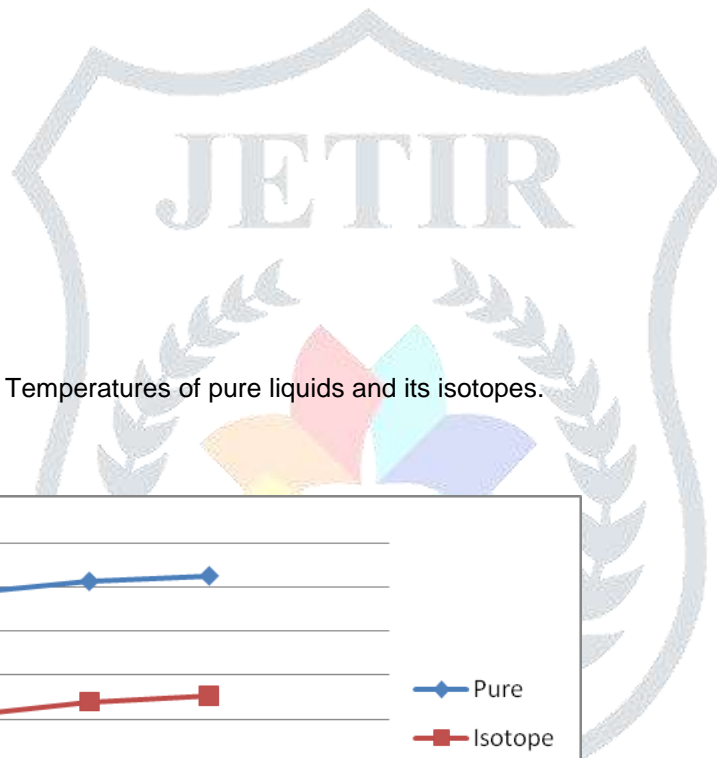
| H ₂ O | ρ (gm/cm ³) | u(m/s) | γ | V | θ_D |
|--------------------------------|------------------------------|--------|----------|--------|------------|
| 293.15 | 0.9991 | 1483 | 1.37 | 18.03 | 157.20 |
| 303.15 | 0.9983 | 1511 | 1.35 | 18.05 | 159.36 |
| 313.15 | 0.9967 | 1531 | 1.33 | 18.08 | 160.65 |
| 323.15 | 0.9940 | 1545 | 1.31 | 18.12 | 161.28 |
| D ₂ O | | | | | |
| 293.15 | 1.0524 | 1386 | 1.35 | 19.03 | 143.48 |
| 303.15 | 1.0509 | 1413 | 1.33 | 19.06 | 145.51 |
| 313.15 | 1.0489 | 1434 | 1.31 | 19.10 | 146.92 |
| 323.15 | 1.0471 | 1448 | 1.29 | 19.13 | 147.62 |
| C ₆ H ₆ | | | | | |
| 293.15 | 0.9376 | 1327 | 1.40 | 83.31 | 85.05 |
| 303.15 | 0.9311 | 1280 | 1.38 | 83.89 | 81.51 |
| 313.15 | 0.9256 | 1232 | 1.36 | 84.39 | 77.98 |
| 323.15 | 0.9199 | 1184 | 1.35 | 84.91 | 74.49 |
| C ₆ D ₆ | | | | | |
| 293.15 | 0.9711 | 1246 | 1.38 | 86.65 | 78.51 |
| 303.15 | 0.9653 | 1201 | 1.37 | 87.17 | 75.21 |
| 313.15 | 0.9594 | 1156 | 1.35 | 87.71 | 71.95 |
| 323.15 | 0.9541 | 1111 | 1.33 | 88.20 | 68.75 |
| C ₆ H ₁₂ | | | | | |
| 293.15 | 0.8823 | 1280 | 1.43 | 95.38 | 78.94 |
| 303.15 | 0.8772 | 1230 | 1.41 | 95.94 | 75.39 |
| 313.15 | 0.8713 | 1180 | 1.39 | 96.59 | 71.87 |
| 323.15 | 0.8661 | 1130 | 1.38 | 97.17 | 68.42 |
| C ₆ D ₁₂ | | | | | |
| 293.15 | 0.9435 | 1153 | 1.40 | 102.01 | 69.02 |
| 303.15 | 0.9377 | 1110 | 1.38 | 102.64 | 66.03 |
| 313.15 | 0.9312 | 1068 | 1.36 | 103.36 | 63.14 |
| 323.15 | 0.9255 | 1026 | 1.35 | 104.00 | 60.29 |
| CH ₃ OH | | | | | |
| 293.15 | 0.8895 | 1122 | 1.42 | 36.02 | 95.64 |
| 303.15 | 0.8848 | 1089 | 1.41 | 36.21 | 92.27 |
| 313.15 | 0.8788 | 1057 | 1.39 | 36.46 | 89.00 |
| 323.15 | 0.8735 | 1024 | 1.37 | 36.68 | 85.71 |
| CH ₃ OD | | | | | |
| 293.15 | 0.8990 | 1103 | 1.42 | 36.76 | 93.28 |
| 303.15 | 0.8943 | 1071 | 1.40 | 36.96 | 90.02 |
| 313.15 | 0.8882 | 1040 | 1.38 | 37.21 | 86.87 |
| 323.15 | 0.8824 | 1008 | 1.37 | 37.45 | 83.69 |

C₂H₅OH

| | | | | | |
|--------|--------|------|------|-------|-------|
| 293.15 | 0.8890 | 1167 | 1.42 | 51.82 | 88.13 |
| 303.15 | 0.8835 | 1133 | 1.41 | 52.15 | 85.02 |
| 313.15 | 0.8788 | 1098 | 1.39 | 52.42 | 81.91 |
| 323.15 | 0.8734 | 1064 | 1.37 | 52.75 | 78.90 |

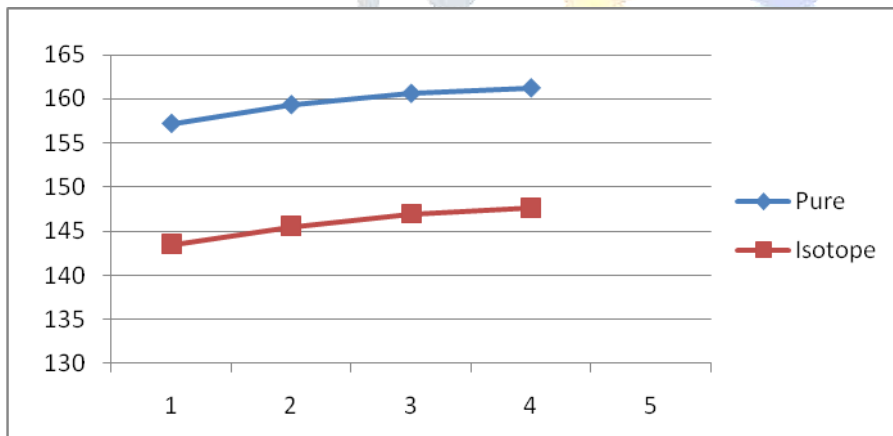
C₂H₅OD

| | | | | | |
|--------|--------|------|------|-------|-------|
| 293.15 | 0.8979 | 1145 | 1.42 | 52.43 | 86.03 |
| 303.15 | 0.8930 | 1111 | 1.40 | 52.72 | 82.97 |
| 313.15 | 0.8877 | 1077 | 1.38 | 53.03 | 79.94 |
| 323.15 | 0.8822 | 1043 | 1.37 | 53.36 | 76.96 |

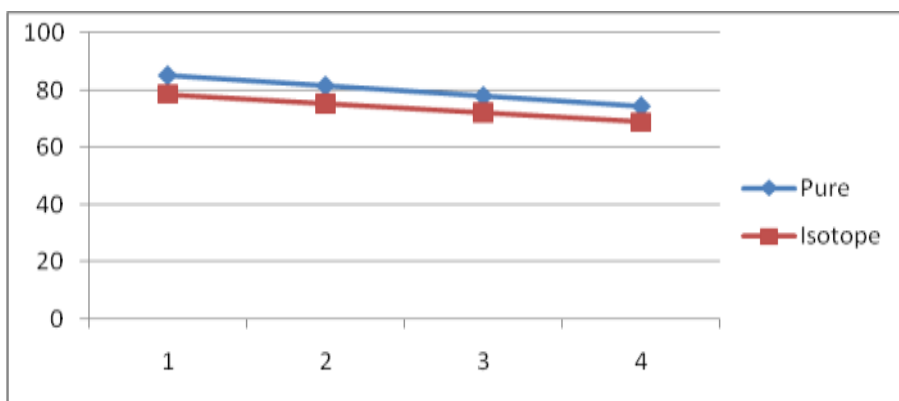


Graph: Variation of Debye Temperatures of pure liquids and its isotopes.

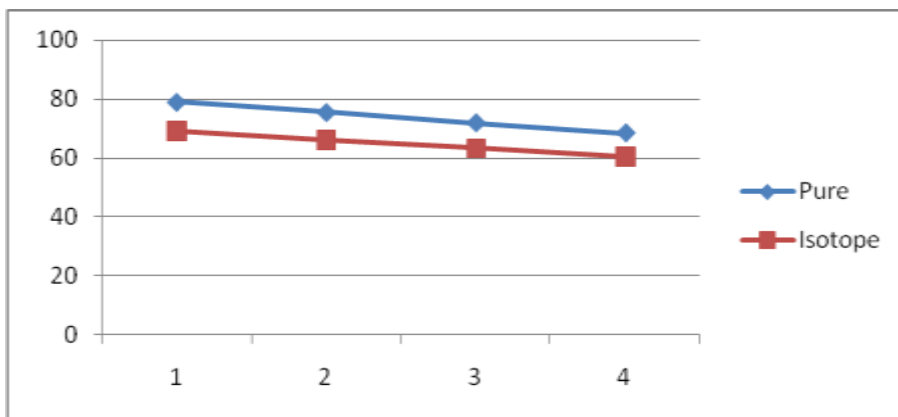
H₂O & D₂O



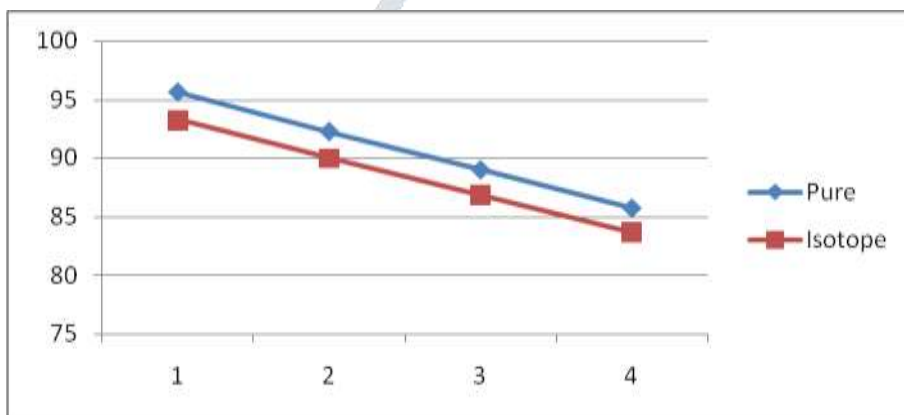
C₆H₆ & C₆D₆



C_6H_{12} & C_6D_{12}



CH_3OH & CH_3OD



C_2H_5OH & C_2H_5OD

