



TWO-DIMENSIONAL FLOW DUE TO THE STRETCHING OF THE SHEET OVER A NATURALLY PERMEABLE BRINKMAN BED

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Abstract:

The study two-dimensional flow of a viscous incompressible fluid through a naturally permeable bed bounded by a stretching sheet has been investigated. The motion being caused solely by stretching of the sheet. The flow behavior due to the stretching of the sheet over a naturally permeable bed using Brinkman's equation, is examined. The results for velocity distribution, stream function, pressure coefficient and coefficient of skin-friction at the stretching sheet have been obtained, drawn for different stretching and permeability parameters, and discussed.

1. Introduction

The fluid flow within a permeable material is represented by a macroscopic law known as Darcy's law. Brinkman(1) further suggested an empirical modification of Darcy's law which was later realized by Tam(2) that the proper way to derive this equation is through the Foldy's(3) approximation. Since then the Brinkman equation has become a major tool in the theoretical investigation of flow in porous media. Many researchers (4-9) also work on stretching sheet. Brinkman has proposed the plausible modification of Darcy's law within effective viscosity considerations. Most applications of the Brinkman equation Saffman (10), Happel(11) assume the effective viscosity and the pure fluid viscosity as same.

Flow in the boundary layer over moving solid surfaces was much investigated. Many researchers (12-17) studied the flow caused by stretching, too. However the flow caused by a stretching sheet in a porous bed does not seem to receive much attention.

Flow and heat transfer effect through porous medium adjacent to a stretching sheet are discussed by Cortell(18), Liu(19) and others. Darcy Brinkmann model is used by Nazar(20) to examine the flow due to stretching sheet in a porous medium. Chauhan and Rastogi(21) studied unsteady flow through a porous medium past a

stretching sheet and heat transfer in the presence of magnetic field. Kiwan and Alnimir (22) investigated similarly solutions for boundary layer flows.

In this paper the steady flow of viscous incompressible through a permeable bed caused by stretching a deformable sheet over the porous matrix, is studied. The flow in the porous matrix is governed by the Brinkman equations. The effect of permeability on the velocity distribution stream function, pressure coefficient and coefficient of skin friction at the stretching sheet have been discussed.

2. FORMULATION OF THE PROBLEM

The two dimensional steady laminar flow of a viscous incompressible fluid through a porous bed is considered. Fluid motion is solely due to the stretching of the sheet over the upper surface of the porous bed filled with fluid. A Cartesian coordinate system is used with the origin at the top of the stretching sheet and y -axis normal to it. The stretching sheet is stretched by introducing two equal and opposite forces so that the position of the point $(0,0)$ remains unchanged.

The flow in the porous matrix $(0 \leq y < -\infty)$ is governed by the Brinkman's equations:

$$v\left(\frac{\partial^2 u}{\partial x^2} + \frac{\partial^2 u}{\partial y^2}\right) - u\frac{v}{K} - \frac{1}{\rho}\frac{\partial p}{\partial x} = 0, \quad (1)$$

$$v\left(\frac{\partial^2 v}{\partial x^2} + \frac{\partial^2 v}{\partial y^2}\right) - v\frac{v}{K} - \frac{1}{\rho}\frac{\partial p}{\partial y} = 0, \quad (2)$$

And

$$\frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} = 0, \quad (3)$$

Where u, v are the fluid velocity components along the x and y directions in the porous matrix ; p the pressure ; ρ the density ; ν the kinematic viscosity; and K the permeability of the porous medium.

The boundary conditions are:

$$\text{At } y = 0, u = Cx, v = 0, \text{ as } y \rightarrow -\infty, u \rightarrow 0, \quad (4)$$

where $u = Cx$ represents the stretching velocity of the sheet with $C > 0$.

3. METHOD OF SOLUTION

Let

$$u = Cxf'(\eta)$$

$$v = -Chf(\eta), \text{ and}$$

$$\eta = \frac{y}{h} \quad (5)$$

where a prime denotes differentiation with respect to η and h is the average pore diameter. Substituting (5) into equation (1) to (4), we have

$$\frac{C^2 x}{R} \left(\frac{1}{\beta} f'(\eta) - f'''(\eta) \right) = -\frac{1}{\rho} \frac{\partial p}{\partial x} \quad (6)$$

$$\frac{C^2 h}{R} \left(f''(\eta) - \frac{1}{\beta} f(\eta) \right) = -\frac{1}{\rho h} \frac{\partial p}{\partial \eta} \quad (7)$$

And the corresponding boundary conditions are

$$\text{at } \eta = 0, f(0) = 0, f'(0) = 1,$$

$$\text{and as } \eta \rightarrow -\infty, f'(-\infty) = 0 \quad (8)$$

$$\text{Where } R = \frac{Ch^2}{\nu} \quad \text{and} \quad \beta = \frac{K}{h^2}.$$

4. SOLUTIONS

The solutions are obtained on solving (6) and (7) under the boundary conditions (8). Thus, we have

$$f'(\eta) = e^{\frac{\eta}{\sqrt{\beta}}} \quad (9)$$

And

$$f(\eta) = \sqrt{\beta} \left(e^{\frac{\eta}{\sqrt{\beta}}} - 1 \right) \quad (10)$$

The stream function for the porous region may be written as

$$\Psi = -R \xi f(\eta), \quad (11)$$

$$\text{where } \xi = \frac{x}{h}$$

The shearing stress on the stretching sheet is

$$\tau = \mu \frac{\partial u}{\partial y} = \mu C x \frac{1}{h} f''(0)$$

Thus the coefficient of skin- friction is given by

$$(C_f)_{\eta=0} = \frac{\tau}{\mu(\nu/h)/h} = -R \xi f''(0) \quad (12)$$

The pressure coefficient is given by

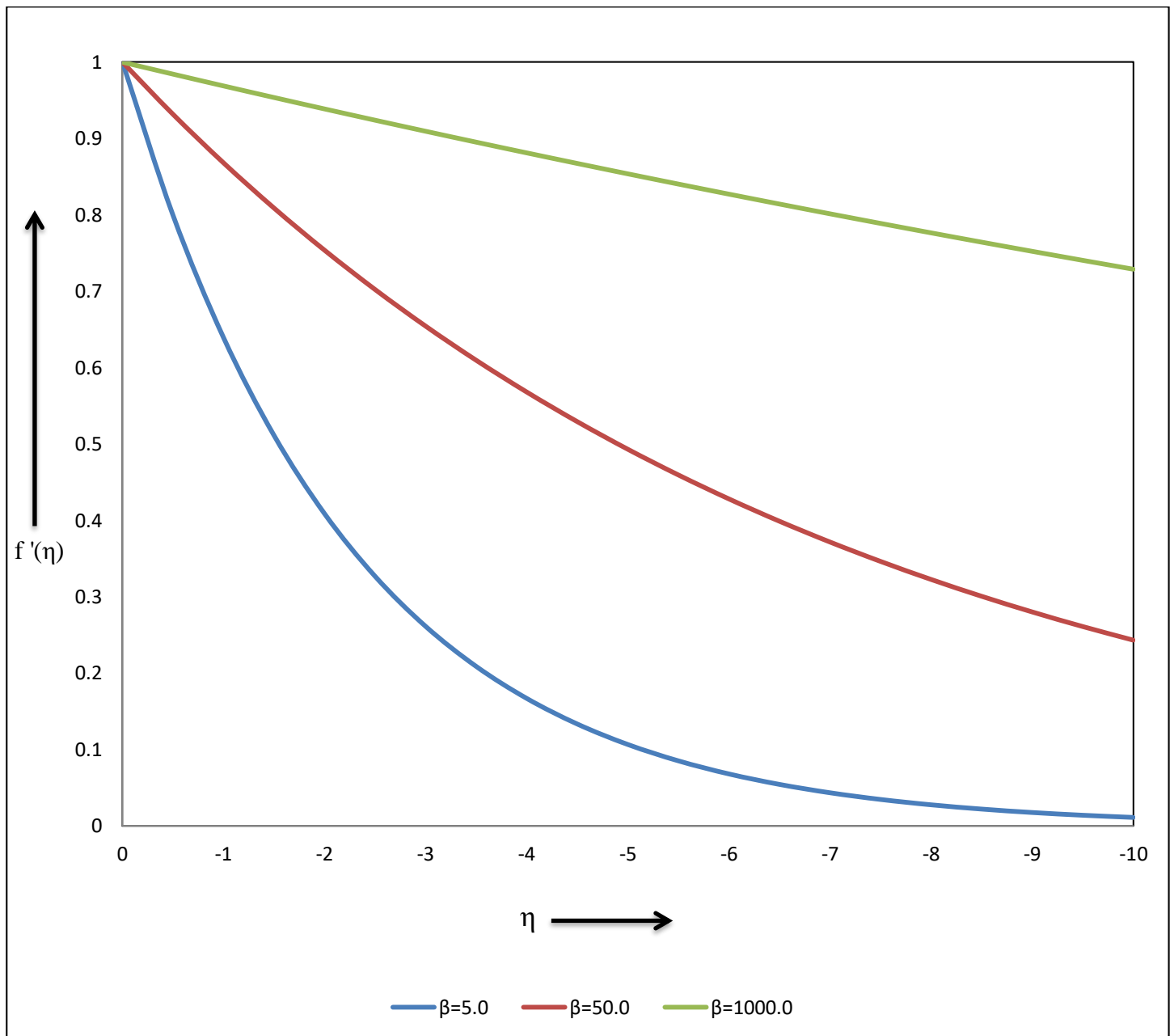
$$C_p = \frac{p}{\rho(v/h)^2} = -\frac{R}{\sqrt{\beta}} \eta$$

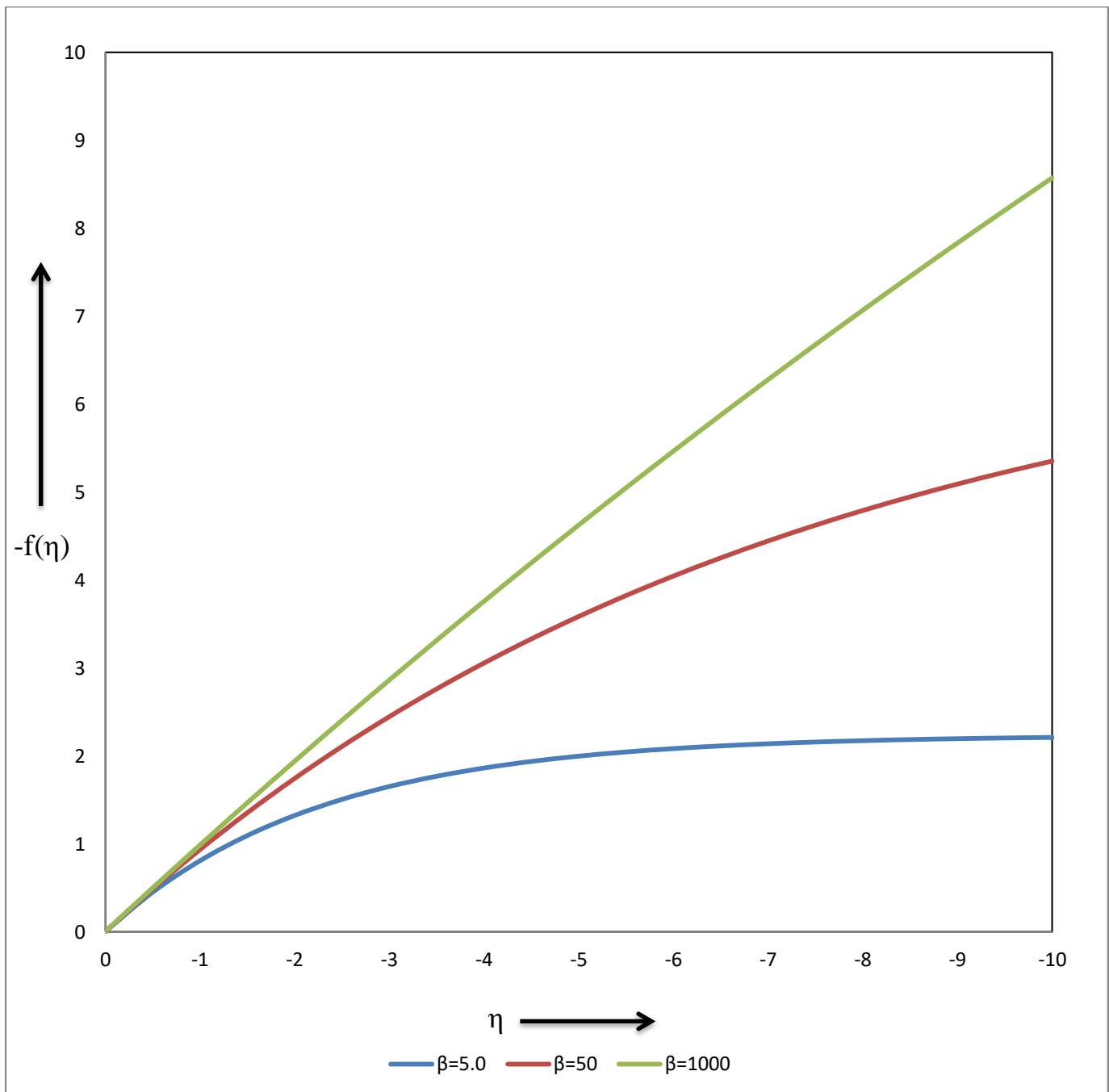
(13)

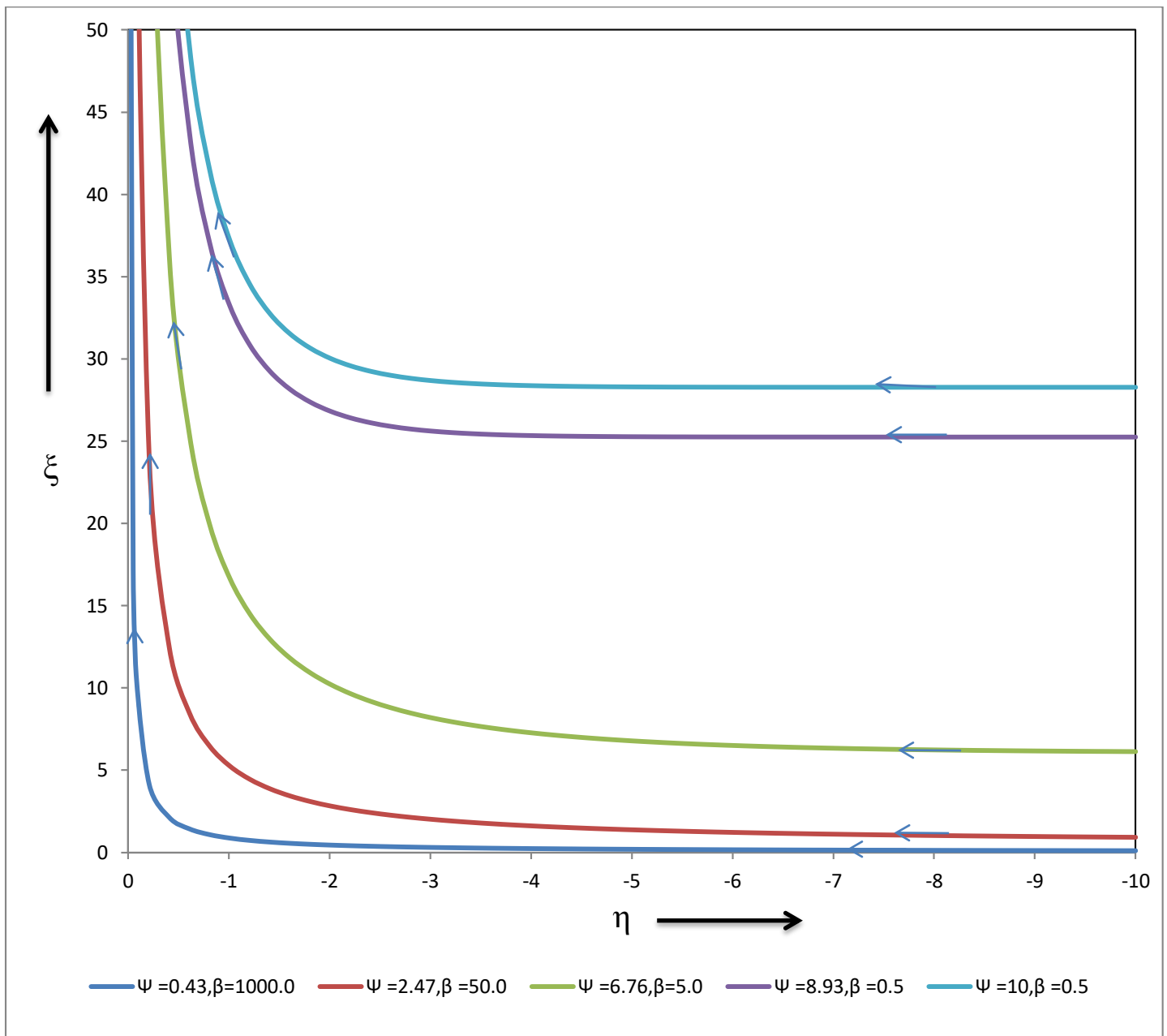
5. DISCUSSION

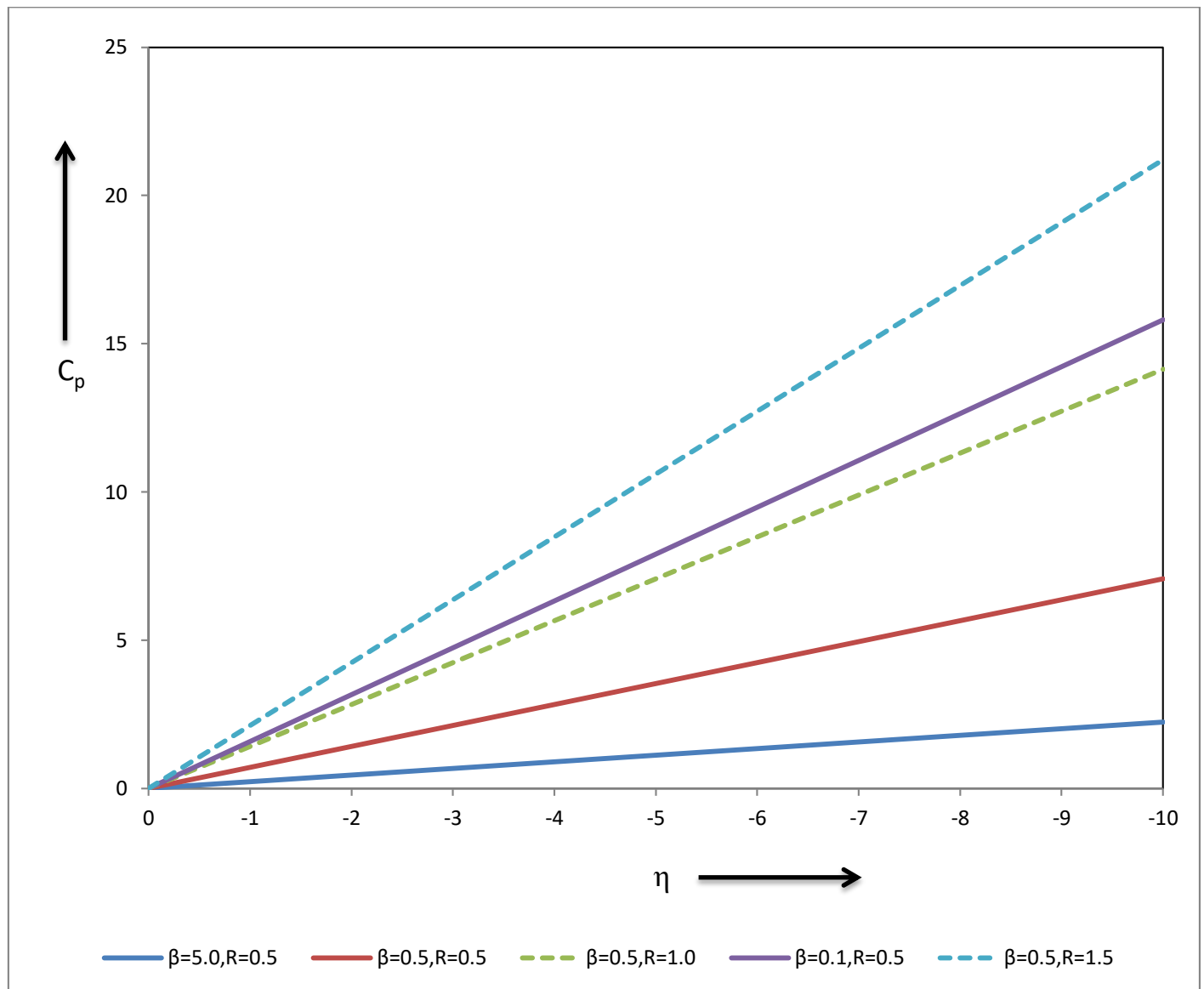
We have drawn velocity profiles and streamlines for the different permeability parameters in Fig 1 - 3. It is found that the fluid velocity components u and v increase with the increase in stretching parameter R and the permeability parameter β . The streamlines show that the fluid is thrownout near the sheet due to stretching and in order to fill the gap the fluid rushes from infinity causing inward flow in the permeable bed. Figures 4 and 5 show the variation of pressure coefficient and the coefficient of skin friction at the stretching sheet for different values of R and β . It is noted that both (C_p and C_f) increase with the increase in R , while both decrease with the increase in β .

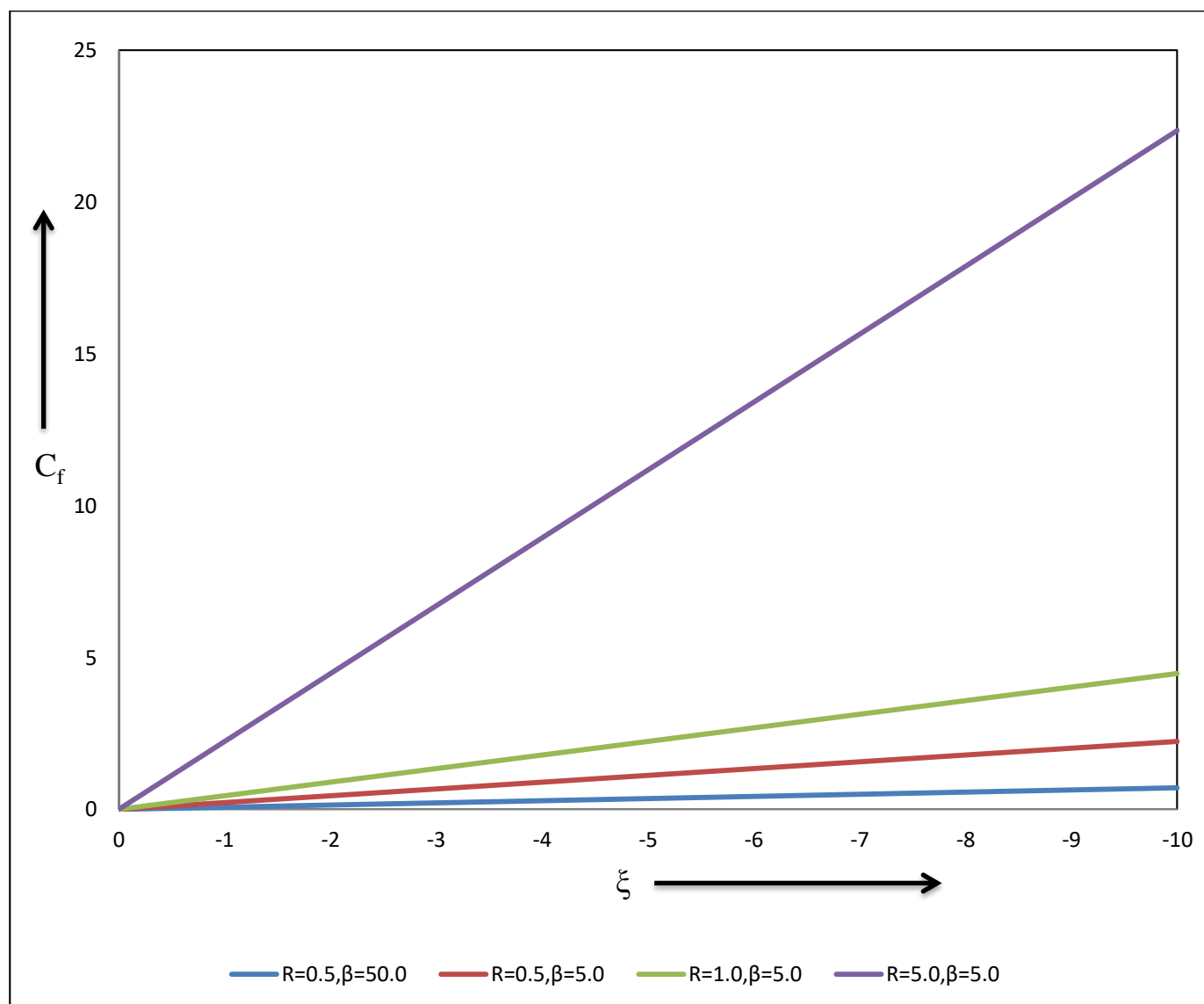


Fig. 1 $f'(\eta)$ Vs η for $R=0.5$

Fig. 2 $-f(\eta)$ Vs η for $R=0.5$

Fig.3 Streamlines Pattern for $R = 0.5$

Fig.4 C_p Vs η

Fig. 5 C_f Vs ξ **REFERENCES:**

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