

Analytical study of wave propagation in micro polar elastic medium

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Abstract

In present investigation, propagation of wave in micropolar elastic medium at non-free surface is discussed. The amplitude ratio's of longitudinal displacement wave (LD wave), coupled transverse displacement (CD-I) and transverse rotational wave (CD-II) are obtained for incident waves.

Key Words: Micropolar elastic medium, non-free surface, wave propagation, amplitude ratio.

1 Introduction

Classical theories of elasticity are not able to examine the behaviour of materials having brous or course grain structure etc. When the microstructure of the material was considered to be rigid, it leads to the micropolar theory. This theory is more dependable for geological materials like solis and rocks as it accounts the intrinsic rotation and estimates the inner structure of the material. Eringen [1] introduced a new formulation of equations in thermoelasticity which was known as the equations for the micropolar elastic theory. Sharma [2] investigated impact of relaxation times and two temperatures on coe cients of re ection in a half-space of micropolar thermoelastic solid.

Fu and Wei [3] investigated the transmission and re ection problem at the imperfect interface of the coupled transverse displacement and transverse rotational waves between two dissimilar micropolar solids. They discussed the impact of imperfect degree of interface on the transmission and the re ection coe cients. Khurana and Tomar [4] observed propagation of plane waves (two longitudinal waves and two sets of coupled transverse waves) for an nonlocal isotropic microp-olar solid and derived re ection coe cients and energy ratios when these waves incidents at stress-free boundary. Singh et.al [5] considered problem on Rayleigh wave for an rotating half-space in an orthotropic micropolar material and solved equations for the surface wave in the half space. They obtained the results to show the in uence of orthotropy, rotation and nondimensional frequency of the Rayleigh wave.

Zhang et.al [6] calculated the amplitude ratios of re ectioned waves for di erent incident waves and also, re ection coe cients in terms of energy ux ratios at non-free surface of a micropolar elastic half-space. Hassanpour and Heppler [7] reviewed the linear isotropic theory of micropolar elasticity with special attention on the notation, which are used for the representation in the micropolar elastic moduli and the experimental actions are taken to measure them. Videla and Atroshchenko [8] derived the analytical solution subjected to a remote uni-axial tension for the problem of a circular micropolar inhomogeneity in an in nite micropolar plate in homogeneous imperfect interface. They showed dependence of stress concentration factors on the micropolar material constants .

Gade and Rangunath [9] explored reduced micropolar theory to replicate ground motion during an earthquake. They calculated the expressions of ground displacement and rotational motions analytically for the case of buried seismic source. Singh [10] investigated a problem on Rayleigh surface

wave in an isotropic micropolar elastic solid half-space with impedance boundary conditions and derived a secular equation for non dimensional speed of the Rayleigh wave, which depends upon various parameters of material, frequency, micro-rotation and impedance parameters. Fan and Cheng [11] presented a elastic model set based on micromechanics in the framework of micropolar theory having two-phase FGMs to study the impact of size on the effective properties of the FGM and compared those results with experimental data.

2 Field Equations

Following Eringen [1], the basic equations and constitutive relations in micropolar elastic medium are:

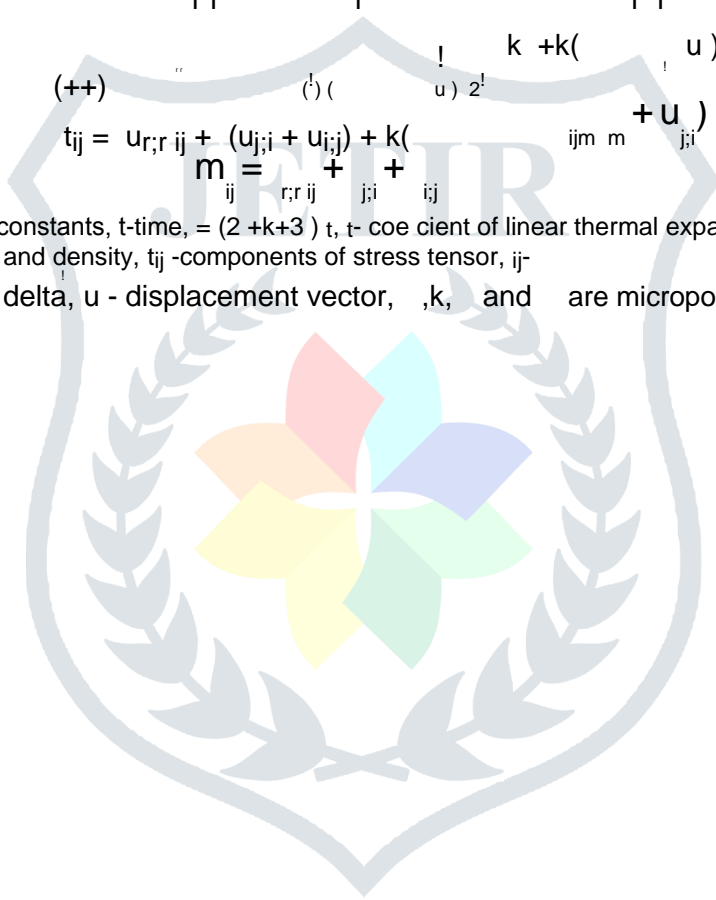
$$\rho \left(\frac{\partial^2 u_i}{\partial t^2} + k_1 \frac{\partial u_i}{\partial t} \right) = \frac{\partial \sigma_{ij}}{\partial x_j} + \frac{\partial \tau_{ij}}{\partial x_j} + \rho \alpha_i \theta \quad (1)$$

$$\rho \left(\frac{\partial^2 \phi_i}{\partial t^2} + k_2 \frac{\partial \phi_i}{\partial t} \right) = \frac{\partial \tau_{ij}}{\partial x_j} + \rho \beta_i \theta \quad (2)$$

$$\sigma_{ij} = \lambda \epsilon_{ij} + \mu (e_{ij} + e_{ji}) + k_3 (\phi_i \phi_j + \phi_j \phi_i) \quad (3)$$

$$\tau_{ij} = \mu (e_{ij} + e_{ji}) + k_4 (\phi_i \phi_j + \phi_j \phi_i) \quad (4)$$

λ, μ - Lamé's constants, t -time, α_i - coefficient of linear thermal expansion, C, ρ - specific heat and density, t_{ij} - components of stress tensor, δ_{ij} - Kronecker delta, u - displacement vector, k_1, k_2 and k_3, k_4 are micropolar constants,



m_{ij} - couple stress tensor components, ϵ_{ijk} is alternating tensor, k is microrotation vectors,.

$$r = i \frac{\partial}{\partial x} + j \frac{\partial}{\partial y} + k \frac{\partial}{\partial z} ; r^2 = \frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2} + \frac{\partial^2}{\partial z^2} :$$

3 Formulation and the solution

We have taken a homogeneous, isotropic with micropolar in elastic half space on non free surface. The rectangular cartesian co-ordinate system (x_1, x_2, x_3) having origin at interface $x_3 = 0$ is considered along with x_3 -axis pointing normally into medium. Plane waves in x_1, x_3 -plane are considered in which wave front is parallel to x_2 -axis, therefore all variables will depend on x_3, x_1 and t . Thus problem considered in two dimensional, so we take

$$\sigma = (0; \tau; 0); \quad u = (u_1; 0; u_3) \tag{5}$$

To ease the solution, quantities having no dimensions are introduced as follows:

$$x_1^0 = \frac{l_1}{c_1} x_1; \quad x_3^0 = \frac{l_1}{c_1} x_3; \quad u_1^0 = \frac{l_1}{c_1} u_1; \quad u_3^0 = \frac{l_1}{c_1} u_3; \\ t_0 = \frac{1}{c_1} t; \quad t_0^0 = \frac{1}{c_1} t; \quad 0 = \frac{2}{k} c_1^2; \quad m_{ij}^0 = \frac{m_{ij} l_1}{c_1} : \tag{6}$$

where

$$c_1^2 = \frac{2 + k}{2} \quad \text{and} \quad l_1 = \frac{k}{j} :$$

The expression related to components of displacement are expressed by using Helmholtz decomposition, therefore u_3 and u_1 are related to the and (scalar potential functions) having no dimensions are given by

$$u_1 = \frac{\partial \phi}{\partial x_1} - \frac{\partial \psi}{\partial x_3} ; \quad u_3 = \frac{\partial \psi}{\partial x_3} + \frac{\partial \phi}{\partial x_1} : \tag{7}$$

using equations (5)-(6) in (1)-(4) and assuming the motion to be harmonic and for solving the equations we assume solutions in the form

$$(\sigma; \tau; \tau) = (0; 0; 0_2) e^{(x_1 \sin \alpha + x_3 \cos \alpha + t)}$$

where denoted as wave number, is known as α , is angle of inclination and quantities such as $0; 0; 0_2$ are arbitrary constants. Using the values of $\sigma; \tau; \tau$ we obtained following equations

$$(c_1^2 + A^2 + B)(c_2^2) = 0 \tag{8}$$

$$(c_2^2 - 1) = 0 \tag{9}$$

where (c_1, c_2, c_3) represents the velocity of various waves; u_1, u_2, u_3 are velocities of the longitudinal displacement (LD) wave, coupled transverse displacement (CD-I) wave and transverse rotational (CD-II) wave respectively and

$$a_1 = \frac{(c_1^2 + k)}{c_1^2}; a_2 = \frac{k}{c_1^2}; a_3 = \frac{k^2}{2c_1^4}; a_4 = \frac{k}{c_1^2}; a_5 = \frac{c_1^2}{c_1^2};$$

$$a_6 = \frac{2k}{c_1^2}; a_7 = \frac{k}{c_1^2}; a_8 = \frac{k}{c_1^2}; a_9 = \frac{k}{c_1^2}; a_{11} = \frac{k^2}{c_1^2};$$

$$a_{12} = \frac{k}{c_1^4}; A = \frac{a_2 a_6}{a_4^2}; B = \frac{a_3 a_5}{a_6^2}; c_3^2 = 1;$$

4 Boundary conditions

Appropriate conditions at surface $x_3=0$ are

$$(i) \quad t_{33} = S_{11}u_3; \tag{10}$$

$$(ii) \quad t_{31} = S_{21}u_1; \tag{11}$$

$$(iii) \quad m_{32} = S_3 u_2 \tag{12}$$

We assume that the values of u_1, u_2, u_3

$$u_1 = A_0 e^{k(x_1 \sin \theta_1 + x_3 \cos \theta_1) + i\omega t} + A_1 e^{k(x_1 \sin \theta_1 + x_3 \cos \theta_1) + i\omega t} \tag{13}$$

$$u_2 = \sum B_0 e^{k(x_1 \sin \theta_1 + x_3 \cos \theta_1) + i\omega t} + B_1 e^{k(x_1 \sin \theta_1 + x_3 \cos \theta_1) + i\omega t} \tag{14}$$

$$u_3 = \sum d_i B_0 e^{k(x_1 \sin \theta_1 + x_3 \cos \theta_1) + i\omega t} + d_i B_1 e^{k(x_1 \sin \theta_1 + x_3 \cos \theta_1) + i\omega t} \tag{15}$$

where

$$d_i = \frac{a_2 k_i^2}{a_3}; \quad (i = 1; 2)$$

where the values of d_i are coupling constants. B_0 are the amplitude of incident coupled transverse displacement (CD-I) and transverse rotational wave (CD-II)

and A_0 is the amplitude of the incident L-D wave (Longitudinal Displacement wave). B_i are the amplitude of the reflected coupled waves i.e transverse rotational and transverse displacement wave and A_1 is the amplitude of the reflected L-D wave (Longitudinal Displacement wave). Using Snell's Law defined as follows

$$\frac{\sin \theta_0}{v_0} = \frac{\sin \theta_1}{v_1} = \frac{\sin \theta_2}{v_2} = \frac{\sin \theta_3}{v_3} \tag{16}$$

where

$$v_1 = v_2 = v_3 = v; \text{ at } x_3 = 0; \tag{17}$$

$\theta_1 > \theta_2 > \theta_3$; for incident LDWave
 $\theta_1 = \theta_2 = \theta_3$; for incident CD IWave
 $\theta_1 > \theta_2 > \theta_3$; for incident CD IIWave
 :

Taking the phase for the reflected waves, one can write the equations (16)-(17)

$$\frac{\cos \theta_j}{v_j} = 2 \frac{v_0}{v_j} \sin^2 \theta_0 \tag{18}$$

Following Schoenberg [12], if we write

$$\frac{\cos \theta_j}{v_j} = \frac{\cos \theta_j^0}{v_j} + \frac{c_j}{2v_0} \quad (j = 1; 2; 3);$$

$$\frac{\cos \theta_j^0}{v_j} = \frac{1}{v_j} \operatorname{Re} \left[\frac{v_0}{v_j} \sin^2 \theta_0 \right]; \quad c_j = 2 \operatorname{Im} \left[\frac{v_0}{v_j} \sin^2 \theta_0 \right];$$

where v_j^0 , called real phase speed and θ_j^0 , known as reflection angle and are given by

$$\frac{v_j^0}{v_j} = \frac{1}{\sin^2 \theta_0 + \operatorname{Re} \left[\frac{v_0}{v_j} \sin^2 \theta_0 \right]}$$

and c_j , is known as attenuation in a depth and equals to $(2v_0)/\lambda$ i.e.wavelength of incident wave

Making use of the equation (7) in the conditions given by (10)-(12) and with the use of equations given by (13)-(15), a homogenous system equations is obtained as follows

$$X \quad a_{ij}Z_j = Y_i; \quad (i; j = 1; 2; 3);$$

$$\begin{aligned}
 a_{1p} &= (a_8 - a_7) \rho \left(\frac{p}{0}\right)^2 \left[\left(\frac{0}{p}\right)^2 \sin^2 \theta\right]^2 \sin \theta - \rho S_1 \frac{p}{0} \sin \theta ; \\
 a_{13} &= \frac{2}{3} a_7 \left(\frac{3}{0}\right)^2 \sin^2 \theta + a_8 \left(\frac{3}{0}\right)^2 \left[\left(\frac{0}{3}\right)^2 \sin^2 \theta\right] + \frac{3}{3} S_1 \left[\left(\frac{0}{3}\right)^2 \sin^2 \theta\right]^2 ; \\
 a_{2p} &= \rho^2 a_9 \left(\frac{p}{0}\right)^2 \left[\left(\frac{0}{p}\right)^2 \sin^2 \theta\right] - \left(\frac{p}{0}\right)^2 \sin^2 \theta - a_{11} \rho \left(\frac{0}{p}\right)^2 \sin^2 \theta - \rho S_2 \left[\left(\frac{0}{p}\right)^2 \sin^2 \theta\right]^2 ; \\
 a_{23} &= \frac{2}{3} (a_9 + 1) \left(\frac{3}{0}\right)^2 \sin \theta \left[\left(\frac{0}{3}\right)^2 \sin^2 \theta\right] - S_2 \frac{3}{0} \sin \theta ; \\
 a_{3p} &= a_p \left[a_{12p} \frac{p}{0} \left[\left(\frac{0}{p}\right)^2 \sin^2 \theta\right]^2 + S_3 \right] ; \quad a_{33} = 0 \quad \text{where } (p = 1; 2)
 \end{aligned}$$

5 Conclusion

In this investigation amplitude ratios are calculated numerically for non-free surface in homogenous isotropic micropolar elastic medium. The amplitude ratios are calculated for incident LD-Wave and Coupled waves, namely coupled transverse rotational wave and transverse displacement wave. The results of the problem can be useful to researcher working in the field of seismology.

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